# 9. A PCILO Study of Conformation and Internal Rotation in Mono-substituted Benzenes 

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#### Abstract

Summary The effect of Kékulé representation and hybrid function of O -atoms in the PCILO-CNDO framework of conformation and internal rotation in mono-substituted benzenes $\mathrm{Ph}-\mathrm{X}\left(\mathrm{X}=\mathrm{NH}_{2}, \mathrm{OH}, \mathrm{OCH}_{3}, \mathrm{CH}_{3}, \mathrm{CHO}, \mathrm{NO}_{2}\right)$ is studied. Three variational criteria for the choice of the appropriate third-order energy, proposed to symmetrize the PCILO results, are critically examined in relation with the height of rotational barrier in these molecules. The study shows that, in all cases, the most stable conformation is qualitatively correct predicted by the PCILO method. Since the barrier to internal rotation in the studied aromatic systems arises predominantly from delocalization effect, it is proposed to employ the arithmetic mean of the third-order energy of the two Kékulé structures. In molecules, in which the third-order energy between the two Kékulé structures is larger than $2 \mathrm{kcal} / \mathrm{mol}$, however, the lower third-order energy representation alone seems to be appropriate. In phenol and anisole the $s p^{3}$-hybridization type of the O -atoms offers better values of rotational barrier, whereas in the $s p^{2}$-type the delocalization is overestimated in the planar conformation.


1. Introduction. - The studies of the relationship between the molecular structure of active drugs and their biological effect are based on the assumption that the specific interaction between the active drug and its receptor can be explained and described in the same terms as those for chemical reactions and interactions [1]. The required informations about molecular reactivity parameters can be obtained -- beside experimental methods - from quantum-chemical calculations [2] [3]. In the field of conformational analysis of polyatomic molecules like drugs the PCILO method (perturbation configuration interaction using localized orbitals) is extensively used because of its speed and recognized successes [4-6]. A general description

[^0]of the main features of this method, first introduced by Diner et al. [7], has been given in a recent review [8].

Through the localization of bonding orbitals in the PCILO framework different chemical formulae are possible for the same configuration. Substituents can break the energetical equivalence between the two Kékulé representations of benzenes and heteroaromatic systems. The same holds for some functional groups (carboxylate, nitro, phosphate etc.). To symmetrize the results within the PCILO method three variational criteria were proposed previously: $a$ ) choice of the representation giving the lower zeroth-order energy ( $E_{0}$-criterion) [7]; $b$ ) choice of the representation giving the lower third-order energy ( $E_{3}$-criterion) [9]; c) arithmetical mean between the two third-order energies ( $\mathrm{A}+\mathrm{B} / 2$-criterion) [9].

Since most of the previously reported applications of the PCILO method are focused on qualitative features, the handling of the cited criteria are not well documented. In 1-acetyl-2-(4-methoxybenzyl)-3-pyrroline ( $=1$-( $N$-Acetyl-3-pyr-rolin-2-yl)methyl-4-methoxybenzol) [10] the energetical difference between the two Kékulé structures was calculated as $0.6 \mathrm{kcal} / \mathrm{mol}$ and in benzylfluoride [11] it amounted to $1.2 \mathrm{kcal} / \mathrm{mol}$. To recover the symmetry in the former molecule the $E_{0}$-criterion and in the latter molecule the ( $\mathrm{A}+\mathrm{B} / 2$ )-criterion was used.

Another study [12] shows that, depending on the cis- or trans-position of the carbonyl group in nicotinamide, benzaldehyde and benzoic acid to the double bond arrangement of the two Kékulé structures, a different potential function of internal rotation was obtained. In a study of phenethylamine [9] it is shown that in the free base of phenethylamine the three criteria give conformational results in satisfactory accord, while in the protonated species the shape of the conformational surface showed some differences. The position of the minima does not change substantially. Compared with ab initio STO-3G calculations the best fit is obtained with the $E_{0}$-criterion. The three versions give even quite different pictures of the rotational surface of zwitter-ionic glycine [13]. Passing from the $E_{0^{-}}$to the $E_{3^{-}}$and then to the $(\mathrm{A}+\mathrm{B} / 2)$ criterion the potential surface becomes gradually smoother, whereas the ( $A+B / 2$ )criterion can grouped together with $a b$ initio STO-3G and CNDO/2 results giving fairly similar conformational maps. In the conformational analysis of benzylideneaniline and azobenzene [14] it seems not reasonable, to consider only one Kékulé structure in the calculation. The arithmetic mean value is in agreement with experimental determinations of the most stable conformation.

Another point of interest in the PCILO method is the hybridization type of O-atom adjoined to aromatic systems. The two possible descriptions of the localized orbitals of the lone-pair electrons of the O -atom in diphenyl ether by the canonical hybrid function $s p^{2}$ or $s p^{3}$ are found to be not equivalent [15].

Since the justification of the choice of one of the cited energy-criteria relies upon the comparison of the computed results with experimental data, we undertook a conformational analysis of monosubstituted benzenes, $\mathrm{Ph}-\mathrm{X}\left(\mathrm{X}=\mathrm{NH}_{2}, \mathrm{OH}\right.$, $\mathrm{OCH}_{3}, \mathrm{CH}_{3}, \mathrm{CHO}, \mathrm{NO}_{2}$ ), for which the most stable conformations as well as barriers to internal rotation are well-known from experiments. Furthermore, these compounds are important constituents of larger molecules of pharmacological interest, often used in drug design for substituent-variated properties of reaction centers.
2. Results and discussion. - All calculations were performed on a $C D C$ computer at the Computer Center ETH-Zürich using the QCPE 221 program [16]. The polarities of bonds were optimized for all conformations. The employed geometries are given in Table 1. The results of the PCILO calculations are listed in Table 2 and compared with those obtained from experimental and other theoretical studies. The Kékulé structures are designated by A or B and their mean value of the third-order energies as $(A+B / 2)$. The canonical hybrid functions for the O -atom $s p^{3}$ and $s p^{2}$ are termed $\sigma, \pi$, respectively. The computational results will be discussed in the following separatly for each molecule, in detail.
a) Toluene. By experimental studies a very low sixfold-barrier of $14 \mathrm{cal} / \mathrm{mol}$ hindering the internal rotation of the methyl group has been found [23] [24]. This type of potential function can be considered as the cancellation of two identical potential functions orthogonal to each other, arising from the rotation of the methyl group with a threefold symmetry against the twofold symmetric frame of the phenyl system [60].

Interestingly, the PCILO calculation confirms this view exactly. Kékulé structures A and B (Fig. 1) show separatly a threefold potential function $180^{\circ}$ out-ofplane with a height of barrier $\mathrm{V}_{3}=1.18 \mathrm{kcal} / \mathrm{mol}$ (Fig. 2). Superposing these two Kékulé structures by the mean value of their third-order energies offers a correct

Table 1. Molecular geometries employed for PCILO calculations


Table 2. Comparison of rotational barriers of mono-substituted benzenes computed by PCILO calculations with those oblained by experimental and other theoretical methods

Molecule \begin{tabular}{l}
Kind of <br>
motion

$\quad$ PCILO (this work) $\quad$

Experimental <br>
data
\end{tabular}$\quad$ Other theoretical methods

| $\mathrm{H}_{2}+\mathrm{H}$ | Rotation $\tau$ | A | $\mathrm{V}_{3}=1.174$ | $\mathrm{~V}_{6}=0.014$ | ab initio STO-3G $<0.006$ [25] |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| $\tau$ | B | $\mathrm{V}_{3}=1.173$ | $[23][24]$ |  |  |
| $\square$ |  | $(\mathrm{A}+\mathrm{B} / 2)$ | $\mathrm{V}_{6}=0.012$ |  |  |


| $\mathrm{O}_{\sim}^{(+1)}$, $\mathrm{O}^{(-1)}$ | Rotation $\tau$ | A | $\mathrm{V}_{2}=3.180$ | $\mathrm{V}_{2}=2.8-3.0$ | MINDO/3 | $6.9\left(\tau=90^{\circ}\right)[28]$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | B | $\mathrm{V}_{2}=2.305$ | (3.0 [26]. | MNDO | $4.9\left(\tau=90^{\circ}\right)[28]$ |
| $\bigcirc$ |  |  | $V_{2}=2.742$ | 2.78 [27]? | $a b$ inilio STO-3G | $5.74\left(\tau=0^{\circ}\right)[25]$ |



Rotation $\tau$

| A | $\mathrm{V}_{2}=1.637$ | $\mathrm{~V}_{2}=4.7-4.9$ |
| :--- | :--- | :--- |
| B | $\mathrm{~V}_{2}=4.045$ | $: 4.66[30]$, |
| $(\mathrm{A}+\mathrm{B} / 2)$ | $\mathrm{V}_{2}=2.841$ | $4.9[31] ;$ |

$\mathrm{CNDO} / 2 \quad 1.1\left(\tau=90^{\circ}\right)[32]$
$\mathrm{CNDO} / 2 \quad 1.2\left(\tau=90^{\circ}\right)[33]$
CNDO/2 (STO-3G
optimized) $\quad 1.1\left(\tau=0^{\circ}\right)$ [33]
PCILO $\quad 2.45\left(\tau=0^{\circ}\right)$ [33]
PCILO $3.30\left(\tau=0^{\circ}\right)[34]$
PCILO Kékulé A $1.29\left(\tau=0^{\circ}\right)$ [13]
Kékulé B $3.30\left(\tau=0^{\circ}\right)$ [13]
ab initio STO-3G $6.60\left(\tau=0^{\circ}\right)$ [25]
$a b$ initio STO-3G $6.56\left(\tau=0^{\circ}\right)$ [33]
ab initio STO-3G $5.84\left(\tau=0^{\circ}\right)$ [33]
(partially optimized)

| $\alpha$ | Barrier to inversion | $A=B$ | 7.303 | 1.3-1.6 | CNDO/2 | 6.4 [39] |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| H ¢ ${ }^{\text {ren }}$ |  |  |  | :1.3 [35]. | CNDO/2 | 7.6 ( $a=54^{\circ}$ ) [40] |
|  |  |  |  | 1.61 [36], | CNDO/2 | $5.97\left(a=54^{\circ}\right)[41]$ |
|  |  |  |  | 1.50 [37]; | INDO | $1.2\left(\alpha=39^{\circ}\right)[40]$ |
|  |  |  |  |  | $a b$ initio |  |
|  |  |  |  | $42^{\circ}$ [35], | STO-3G | $2.7\left(a=48^{\circ}\right)[25]$ |
|  |  |  |  | $46^{\circ}$ [36]! | STO-3G | $4.3\left(\alpha=51^{\circ}\right)[42]$ |
|  |  |  |  |  | STO- $(1+1+1) \mathrm{G}$ | $1.1\left(a=38^{\circ}\right)[42]$ |
|  |  |  |  |  | [5.2/2]DZ | $0.9\left(\alpha=39^{\circ}\right)[42]$ |
|  | Rotation $\tau$ | $a=0^{\circ}$ |  | 3.54 [38] | CNDO/2 $2=0^{\circ}$ | 16.06 [41] |
|  |  | A | 11.47 |  | $a=54^{\circ}$ | 4.30 [41] |
|  |  | B | 11.47 |  | INDO $\quad a=0^{\circ}$ | 10.3 [40] |
|  |  | $(\mathrm{A}+\mathrm{B} / 2)$ | 11.47 |  | $a=39.4{ }^{\circ}$ | 7.5 [40] |
|  |  | $a=55^{\circ}$ |  |  | $a b$ initio STO-3G |  |
|  |  | A | 5.816 |  | $\alpha=0^{\circ}$ | 10.7 [43] |
|  |  | B | 7.929 |  |  |  |
|  |  | $(\mathrm{A}+\mathrm{B} / 2)$ | 6.873 |  |  |  |
|  | Rotation $\tau$ | $\sigma \mathrm{A}$ | 3.305 |  |  | $2.31\left(\tau=0^{\circ}\right)[48]$ |
|  |  | $\sigma \mathrm{B}$ | 2.814 | \{3.29 [44]. | CNDO/2 | $2.76\left(\tau=0^{\circ}\right)\left[49{ }^{\circ}\right.$ |
|  |  | $\sigma(\mathrm{A}+\mathrm{B} / 2)$ | 3.060 | 3.26 [45], | $a b$ initio STO-3G | $5.15\left(\tau=0^{\circ}\right)[25]$ |
|  |  | $\pi \mathrm{A}$ | 5.267 | 3.36 [19], | STO-3G | $4.71\left(\tau=0^{\circ}\right)[50]$ |
|  |  | $\pi \mathrm{B}$ | $4.612$ | 3.47 [46], | STO-3G | $4.08\left(\tau=0^{\circ}\right)[51]$ |
|  |  | $\pi(A+B / 2)$ | 4.940 | 3.56 [471) | (INDO-optimized) |  |


| Molecule | Kind of motion | PCILO (this work) |  | Experimental data | Other theoretical methods |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Rotation $\tau_{1}$ | $\sigma \mathrm{A}$ | 2.751 | not certain | CNDO/2 | $1.76\left(\tau_{1}=0^{\circ}\right)$ [54] |
|  |  | $\sigma \mathrm{B}$ | 1.724 | \{3.61 [52], | CNDO/2 | $1.00\left(\tau_{1}=0^{\circ}\right)$ [49] |
|  |  | $\sigma(\mathrm{A}+\mathrm{B} / 2)$ | 2.238 | $>1.1[53]$, | INDO | $\left(\tau_{1}=90^{\circ}\right)[50]$ |
|  |  | $\pi \mathrm{A}$ | 5.030 |  | MINDO/3 | $1.28\left(\tau_{1}=0^{\circ}\right)$ [55] |
|  |  | $\pi \mathrm{B}$ | 3.550 |  | $a b$ initio STO-3G | $0.06\left(\tau_{1}=0^{\circ}\right)[25$, |
|  |  | $\pi(A+B / 2)$ | 4.290 |  |  | $51,56]$ |
|  |  |  |  |  | $a b$ initio STO-3G | $0.94\left(\tau_{1}=0^{\circ}\right)[50]$ |
|  |  |  |  |  | $a b$ initio STO-3G | $1.34\left(\tau_{1}=0^{\circ}\right)[50]$ |
|  | Rotation $\tau_{2}$ | with $\tau_{1}=0^{\circ}$ |  | 1.8-2.8 | $\operatorname{INDO} \tau_{1}=20^{\circ}$ | 4.02 [59] |
|  |  | $\sigma \mathrm{A}$ | 3.403 | !1.79 [57], | $\tau_{1}=90^{\circ}$ | 1.05 [59] |
|  |  | $\sigma \mathrm{B}$ | 3.301 | 1.79 [58], |  |  |
|  |  | $\sigma(\mathrm{A}+\mathrm{B} / 2)$ | 3.352 | 2.8 [59]\} |  |  |
|  |  | $\pi \mathrm{A}$ | 3.600 |  |  |  |
|  |  | $\pi \mathrm{B}$ | 3.201 |  |  |  |
|  |  | $\pi(A+B / 2)$ | 3.401 |  |  |  |



Kékulé A


Kékulé B

Fig. 1. Designation of the Kékulé structures of toluene


Fig. 2. Rotation of the methyl group in toluene
sixfold potential function with a barrier of $\mathrm{V}_{6}=12 \mathrm{cal} / \mathrm{mol}$. The most stable conformation is found to be that one with one H -atom in the plane of the phenyl system.
b) Nitrobenzene. Supported by the experimentally determined planar structure of nitrobenzene in its ground state [26] [29] it can be assumed that the N -atom enters into chemical bonding in an $s p^{2}$-hybridized state. Three of the $s p^{2}$-orbitals form planar $\sigma$-bonds with the O -atoms and the ring C -atom. The remaining N -$p$-orbital forms a bond with one of the O -atoms, in which the ring electrons are largely delocalized (Fig. 3).

Assuming a symmetric nitro group the Kékulé structures $\mathrm{A}, \mathrm{A}_{1}, \mathrm{~B}$ and $\mathrm{B}_{1}$ can be reduced to two different arrangements with respect to the substituent ( $A=B_{1}$ and $\mathrm{B}=\mathrm{A}_{1}$ ). Therefore, the arithmetical mean of the two Kékule structures seems to describe the system a priori in an appropriate manner.


Fig. 3. Designation of the Kékulé structures of nitrobenzene

In the first definition the negative-charged O -atom is considered to be $s p$ hybridized. Within this description an unrealistically large dipole moment of 41.2 D is calculated compared to the experimental value of 4.22 D [61]. Evident by the poor convergence of the CI-terms (Table 3), especially the very large delocalized single excitation, the degree of delocalization in the system becomes too large to be treated as perturbation in the localized picture of the PCILO framework [62].

In the second definition the negative-charged O -atom is described as $s p^{3}$ hybridized, which leads to a more reasonable dipole moment of 5.55 D and a better convergence in the CI-terms. In spite of the ambiguity of the $s p^{3}$-hybridization of the negative charged O -atom the conformational results of the PCILO calculations are in good agreement with experiment (Table 2). In both Kékulé structures the planar conformation is found to be the most stable one and a twofold barrier is hindering the internal rotation of the nitro group. The third-order energies of the two Kékulé structures in the planar conformation are different by $0.88 \mathrm{kcal} / \mathrm{mol}$, favouring the Kékulé structure A. In the orthogonal conformation the two representations are equivalent. Both the $E_{0^{-}}$and the $E_{3}$-criteria lead to the Kékulé structure A as the lower one in energy. Compared with the experimentally determined rotational barrier of $2.8-3.0 \mathrm{kcal} / \mathrm{mol}$ [26] [27] the Kékulé structure A amounts to a better predicted height of barrier ( $3.18 \mathrm{kcal} / \mathrm{mol}$ ) than the Kékulé structure $\mathrm{B}(2.31 \mathrm{kcal} / \mathrm{mol})$. The best agreement with experiment is, however, obtainable from the arithmetical mean value of the two representations ( $2.74 \mathrm{kcal} / \mathrm{mol}$ ) (Table 2).

Table 3. PCILO energy contributions of nitrobenzene ( $\mathrm{kcal} / \mathrm{mol}$ )

| Dihedral angle $\tau$ | $\tau=0^{\circ}$ |  |  | $\tau=90^{\circ}$ |
| :---: | :---: | :---: | :---: | :---: |
| Hybridization type of O-atom | $s p$ | $s p^{3}$ | $s p^{3}$ | $s p^{3}$ |
| Kékulé structure | A | A | B | $\mathrm{A}=\mathrm{B}$ |
| Nuclear repulsion energy | 131939.47 | 131939.47 | 131939.47 | 131416.46 |
| $E_{0}$ (zeroth-order energy) | - 59002.281 | - 59076.882 | -59075.768 | -59076.298 |
| Delocalization energy | -641.406 | -367.990 | -368.521 | - 364.493 |
| Intra-bond correlation energy | - 127.626 | - 127.123 | - 127.325 | - 127.287 |
| Inter-bond correlation energy | - 169.660 | -169.68? | - 171.750 | - 169.423 |
| $E_{2}$ (second-order energy correction) | -938.692 | -664.795 | -667.596 | -661.203 |
| $E_{0}+E_{2}$ | - 59940.973 | - 59741.677 | -59743.364 | -59737.501 |
| Delocal. 2 bonds correl. interaction ${ }^{\text {a }}$ ) | - 11.333 | - 11.344 | - 12.066 | - 10.924 |
| 1 Bond correl. 2 bond correl. interaction | 76.805 | 76.904 | 78.556 | 77.135 |
| Delocal. delocal. interaction | -35.176 | -74.291 | -76.728 | - 76.797 |
| 2 Bond correl. 2 bond correl. interaction | 52.925 | 52.953 | 57.023 | 53.812 |
| $E_{3}$ (third-order energy correction) | 83.221 | 44.222 | 46.785 | 43.226 |
| $E_{0}+E_{2}+E_{3}$ | - 59857.752 | - 59697.455 | -59696.579 | - 59694.275 |
| ${ }^{\text {a }}$ ) Delocal. stands for 'delocalization' and correl. for 'correlation'. |  |  |  |  |



Fig. 4. Designation of the Kékule structures of benzaldehyde

This PCILO result may be contrasted to previous MINDO/3 and MNDO results, which predict the orthogonal conformation to be the most stable conformation [28]. As for the barrier height, an ab initio calculation with STO-3G basis set [25] gave $5.44 \mathrm{kcal} / \mathrm{mol}$, overestimating the experimental value appreciably.
c) Benzaldehyde. Benzaldehyde (Fig. 4) is known to be planar in the ground state and has a twofold barrier with a height of $4.7-4.9 \mathrm{kcal} / \mathrm{mol}$ [30] [31]. This planar conformation is also confirmed from previous and the present PCILO studies as the most stable one (Table 2), though the calculated barrier seems to be sensitive to the employed geometry.

In the present study, the variational criteria for the choice of the appropriate energy value lead to different Kékulé viructures (Table 4) ( $E_{0}$-criterion $\rightarrow$ Kékulé A, $E_{3}$-criterion $\rightarrow$ Kékulé B). A comparison between the values of the two Kékulé structures and experimental data shows that the rotational barrier obtained from Kékulé structure B gives better agreement with experiment (Table 2). In spite of a large energy difference between both Kékulé structures in the planar conformation $\left(\Delta E_{3}^{\mathrm{AB}}=2.4 \mathrm{kcal} / \mathrm{mol}\right)$, it may be concluded that, in the localized picture of the PCILO method, both Kékulé structures have a different weight in the planar conformation, the Kékulé structure B being the dominant form. In view of this, it seems not advisable to choose the ( $\mathrm{A}+\mathrm{B} / 2$ )-criterion for the description of the internal rotation in benzaldehyde. In the orthogonal state the two Kékulé structures are found to be equivalent, as expected.

Compared with other theoretical studies on benzaldehyde, the PCILO method gives a better qualitative picture of the rotational behaviour. The failure of the

Table 4. PCILO energy contributions of benzaldehyde ( $\mathrm{kcal} / \mathrm{mol}$ )

| Dihedral angle $\tau$ | $\tau=0^{\circ}$ |  | $\tau=90^{\circ}$ |
| :---: | :---: | :---: | :---: |
| Kékulé structure | $\overline{\mathrm{A}}$ | B | $\mathrm{A}=\mathrm{B}$ |
| Nuclear repulsion energy | 101959.09 | 101959.09 | 101675.40 |
| $E_{0}$ (zeroth-order energy) | -45105.086 | -45104.958 | -45105.515 |
| Delocalization energy | -312.503 | -313.653 | -309.032 |
| Intra-bond correlation energy | - 115.082 | - 115.131 | - 115.134 |
| Inter-bond correlation energy | - 150.998 | - 149.808 | - 149.092 |
| $E_{2}$ (second-order energy correction) | - 578.582 | - 578.592 | - 573.258 |
| $E_{0}+E_{2}$ | -45683.668 | -45683.550 | -45678.773 |
| Delocal. 2 bonds correl. interaction ${ }^{\text {a }}$ ) | - 12.651 | - 11.821 | -11.415 |
| 1 Bond correl. 2 bond correl. interaction | 62.156 | 61.361 | 61.221 |
| Delocal. delocal. interaction | -44.617 | -44.374 | -45.670 |
| 2 Bond correl. 2 bond correl. interaction | 47.364 | 44.561 | 44.860 |
| $E_{3}$ (third-order energy correction) | 52.252 | 49.724 | 48.993 |
| $E_{0}+E_{2}+E_{3}$ | -45631.418 | -45633.826 | -45629.781 |

${ }^{\text {a }}$ ) Delocal. stands for 'delocalization' and correl. for 'correlation'.

CNDO/2 method using standard geometry is well-known [32] [33] [63]. The height of barrier is also better described by the PCILO method than by the ab initio STO-3G calculation using standard geometry (Table 2).
d) Aniline. For a study of the inversion of the amino group in aniline a pure pyramidal hybridization at the N -atom is assumed (Fig. 5). The angle of pyramidalization $\alpha$ is interrelated to the bond angle $\Varangle \mathrm{HNH}(\beta)$ by [42]

$$
\operatorname{tg} a=\cos (\beta / 2) \cdot \sqrt{\operatorname{tg}^{2} \beta-\operatorname{tg}^{2}(\beta / 2)}
$$

The dihedral angle $\phi$ is given by [42]

$$
\operatorname{tg} \phi=\operatorname{ctg}(\beta / 2) \cdot \sin \alpha
$$

The lone-pair F (Fig. 5) is assumed to be orthogonal to the plane of the phenyl system and the angle $\Varangle \mathrm{CNF}$ is fixed at $90^{\circ}$. In this model the inversion is only a function of the angle of pyramidalization $a$. This angle is varied in $15^{\circ}$ steps within the range of $0^{\circ}$ and $75^{\circ}$.

The PCILO computation correctly predicts the pyramidal structure to be the most stable one with $\alpha=55^{\circ}$. A double-minimum potential function for the inversion of the amino group is found with a barrier of $7.3 \mathrm{kcal} / \mathrm{mol}$. The two Kékulé structures A and B (Fig. 6) have the same zeroth-order energy and the CI-terms give the same third-order energy. The three different criteria are equivalent. The barrier of inversion is found as too high compared with experimental data (Table 2), which is similar to the CNDO/2 results. $A b$ initio STO-3G calculations predicted also a too high barrier to inversion, only an extended basis set yield a quantitatively correct barrier and angle of pyramidalization [42]. On the other hand surprisingly good results were obtained by the INDO method [40]. In the present PCILO calculation the zeroth-order wave function overestimates the delocalization contributions in the CI-perturbative treatment. A similar shortcoming of both PCILO and CNDO/2 is also observed by Weller \& Lochmann [64] in a series of aliphatic amines in predicting the inversion barriers.

By rotating the amino group the two Kékulé structures are no more equivalent (Table 5). The conformational energy map shows (Fig. 7) that the choice of the Kékulé structure in the region for $a$ larger than $50^{\circ}$ is important. With $a=55^{\circ}$ the




Kèkulé B


Kèkulé A


Fig. 5. Definition of angles in aniline

Fig. 6. Designation of the Kékulé structures of aniline


Fig. 7. Conformational energy maps of aniline a) Kékulé structure $A, b$ ) Kékulé structure $B$ (the isoenergy curves run from 0 to $19 \mathrm{kca} / \mathrm{mol}$ in $1 \mathrm{kcal} / \mathrm{mol}$ steps)
rotational barrier for the amino group is calculated for Kékulé structure A as 5.80 $\mathrm{kcal} / \mathrm{mol}$ and for Kékulé structure B as $7.93 \mathrm{kcal} / \mathrm{mol}$. Compared with experimental data Kékulé structure A seems to describe the barrier more reasonably (Table 2). This Kékulé structure is predicted by the $E_{3}$-criterion, whereas on the base of the $E_{0}$-criterion the structure B may be chosen. The pronounced energy difference between the two Kékulé structures in the third-order energy ( $\Delta E_{3}^{\mathrm{AB}}=2.1 \mathrm{kcal} / \mathrm{mol}$ ) during rotation shows similar behaviour as in benzaldehyde. The Kékulé structure A is predominantly stabilized by delocalization contributions in the second- and third-order terms corrected by interbond correlation energy (Table 5). Therefore, the ( $\mathrm{A}+\mathrm{B} / 2$ )-criterion

Table 5. PCILO energy contributions of aniline ( $\mathrm{kcal} / \mathrm{mol}$ )

| Angle of pyramidalization $a$ | $a=0^{\circ}$ | $a=55^{\circ}$ | $a=55^{\circ}$ |  |
| :---: | :---: | :---: | :---: | :---: |
| Dihedral angle $\tau$ | $\tau=0^{\circ}$ | $\tau=0^{\circ}$ | $\tau=90^{\circ}$ |  |
| Kékulé structure | $\mathrm{A}=\mathrm{B}$ | $\mathrm{A}=\mathrm{B}$ | A | B |
| Nuclear repulsion energy | 87413.461 | 87661.348 | 87607.836 | 87607.836 |
| $E_{0}$ (zeroth-order energy) | -36987.945 | -36979.501 | -36980.206 | -36980.235 |
| Delocalization energy | -288.659 | - 302.032 | -294.359 | - 292.488 |
| Intra-bond correlation energy | -101.136 | - 101.474 | - 101.475 | - 101.465 |
| Inter-bond correlation energy | -118.106 | -118.453 | -117.356 | - 118.348 |
| $E_{2}$ (second-order energy correction) | -507.902 | -521.958 | -513.188 | -512.301 |
| $E_{0}+E_{2}$ | -37495.846 | -37501.459 | -37493.394 | -37492.536 |
| Delocal. 2 bonds correl. interaction ${ }^{\text {a }}$ ) | $-10.729$ | - 10.679 | $-10.902$ | - 11.015 |
| 1 Bond correl. 2 bond correl. interaction | 46.220 | 56.327 | 45.895 | 46.351 |
| Delocal. delocal. interaction | -43.463 | -45.212 | -46.031 | -46.250 |
| 2 Bond correl. 2 bond correl. interaction | 36.353 | 36.253 | 35.474 | 36.600 |
| $E_{3}$ (third-order energy correction) | 28.374 | 26.684 | 24.435 | 25.688 |
| $E_{0}+E_{2}+E_{3}$ | -37467.473 | -37474.776 | -37468.960 | -37466.848 |

[^1]would lead to an overestimation of Kékulé structure B and to a too high rotational barrier of $6.87 \mathrm{kcal} / \mathrm{mol}$ compared with experimental data. The PCILO result of the rotation of the amino group is comparable to the INDO and ab initio results, whereas CNDO/2 predicted a too high barrier to rotation with $a=0^{\circ}$ (Table 2).
e) Phenol. It is well-established that phenol is a planar molecule with a rotational barrier of $3.4 \mathrm{kcal} / \mathrm{mol}$ (Table 2). Because the height of barrier may be associated with the moderably large degree of double-bond character of the $\mathrm{C}, \mathrm{O}$-bond, it is interesting to note the effect of the lone-pair description of the O -atom by the canonical hybrid function $s p^{3}$ or $s p^{2}$ ( $\sigma$ - or $\pi$-form). The employed geometries of the fictitious atoms $\mathrm{F}_{\mathrm{i}}$ are defined in Figure 9. In general, we deduce the fictitious bond angle $\Varangle \mathrm{R}_{2} \mathrm{OF}(\phi)$ as a function of the bond angle $\Varangle \mathrm{R}_{2} \mathrm{OR}(a)$ and the dihedral angle $\mathrm{R}_{1} \mathrm{R}_{2} \mathrm{OF}(\gamma)$ for the $\sigma$-form (Fig. 9):
$$
\cos \phi=-\cos (\alpha / 2) \cdot \cos (\gamma / 2)
$$

Assuming $\gamma=109.5^{\circ}$, it follows:

$$
\cos \phi=-0.57715 \cos (\alpha / 2)
$$

For the fictitious bond angle $\Varangle \mathrm{R}_{2} \mathrm{OF}(\phi)$ in the $\pi$-form (Fig. 9) we take $\phi=180^{\circ}-(\alpha / 2)$. The bond angle $\Varangle \mathrm{CCO}$ is taken as $122.2^{\circ}$ in the planar conformation and as $120.0^{\circ}$ in the orthogonal one.

The energy difference of the $\sigma, \pi$-forms and the Kékulé structures during the rotation of the hydroxyl group are shown in Table 6. In all cases phenol is predicted to be planar by PCILO calculation.

The $\sigma$ - and $\pi$-descriptions are energetically not equivalent. Both Kékulé structures of the $\sigma$-form have a better zeroth-order energy ( $\Delta E^{\sigma \pi}=9.5 \mathrm{kcal} / \mathrm{mol}$ ) in the planar and orthogonal conformation. The third-order energy of the $\pi$-form is favored in the planar conformation due to a larger delocalization energy $\left(\Delta E_{3}^{\sigma \pi}=1.1 \mathrm{kcal} / \mathrm{mol}\right)$, whereas in the orthogonal state the $\pi$-form is smaller in energy ( $\Delta E_{3}^{\sigma \pi}=0.8 \mathrm{kcal} / \mathrm{mol}$ ). This is caused mainly from changes in the delocalization energy due to rotation, which is $2.8 \mathrm{kcal} / \mathrm{mol}$ in the $\sigma$-form and $8.2 \mathrm{kcal} / \mathrm{mol}$ for the $\pi$-form, respectively.

In the planar conformation the energy differences between the Kékulé structures of the same hybridization type are $\Delta E_{3 \sigma}^{\mathrm{AB}}=0.5$ and $\Delta E_{3 \pi}^{\mathrm{AB}}=0.66 \mathrm{kcal} / \mathrm{mol}$,


Fig. 8. Designation of the Kékulé structures of phenol
Kékulé A
Kékulé B
Kékulé C


Fig. 9. Geometry of $O$-atom in phenol

Table 6. PCILO energy contributions of phenol ( $\mathrm{kcal} / \mathrm{mol}$ )

| Dihedral angle $\tau$ <br> Hybridization type of oxygen atom Kékulé structure | $\tau=0^{\circ}$ |  |  |  | $\tau=90^{\circ}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\sigma$ |  |  |  |  |  |
|  | A | B | A | B | $\mathrm{A}=\mathrm{B}$ | $A=B$ |
| Nuclear repulsion energy | 87179.100 | 87179.100 | 87179.100 | 87179.100 | 87137.842 | 87137.842 |
| $E_{0}$ (zeroth-order energy) | -40761.781 | -40761.759 | -40751.052 | -40751.034 | -40762.600 | -40752.125 |
| Delocalization energy | -276.062 | - 275.263 | -291.402 | -290.221 | -272.839 | -282.645 |
| Intra-bond correlation energy | -98.191 | -98.199 | -97.578 | -97.584 | -98.195 | -97.580 |
| Inter-bond correlation energy | - 117.146 | - 116.218 | -116.839 | - 115.933 | -116.146 | -115.861 |
| $E_{2}$ (second-order energy correction) | -491.399 | -489.680 | - 505.820 | -503.738 | -487.181 | -496.086 |
| $E_{0}+E_{2}$ | -41253.180 | -41251.439 | -41256.872 | -41254.772 | $-41249.780$ | -41248.21I |
| Delocal. 2 bonds correl. interaction ${ }^{\text {a }}$ ) | - 10.555 | $-10.556$ | - 10.515 | -10.516 | -10.718 | - 10.648 |
| 1 Bond correl. 2 bond correl. interaction | 46.090 | 45.657 | 45.914 | 45.495 | 45.690 | 45.529 |
| Delocal. delocal. interaction | -47.740 | -47.662 | -44.992 | -45.142 | -46.601 | -47.213 |
| 2 Bond correl. 2 bond correl. interaction | 35.502 | 34.607 | 35.384 | 34.506 | 34.828 | 34.724 |
| $E_{3}$ (third-order energy correction) | 23.294 | 22.044 | 25.786 | 24.340 | 23.200 | 22.392 |
| $E_{0}+E_{2}+E_{3}$ | -41229.886 | $-41229.395$ | -41231.086 | -41230.431 | -41226.581 | -41225.819 |

${ }^{\text {a }}$ ) Delocal. stands for 'delocalization' and correl. for 'correlation'.
favoring in both hybrid forms the Kékulé structure A. In the orthogonal conformation the energies of the Kékulé structures of the same hybridization type are degenerate.

The criteria for choosing an appropriate form ( $\sigma$ and $\pi$ ) and Kékulé structure (A and B) are suggested as follows,

$$
\begin{aligned}
& E_{0} \text {-criterion } \rightarrow \sigma \text {-form, } \text { Kékulé structure } \mathrm{A} \\
& E_{3} \text {-criterion } \rightarrow \text { planar: } \pi \text {-form, Kékulé structure } \mathrm{A} \\
& \text { orthogonal: } \sigma \text {-form, Kékulé structure } \mathrm{A}=\mathrm{B}
\end{aligned}
$$

A comparison between experimental data and the results of PCILO calculations shows (Table 2) that the barrier to internal rotation is predicted to be too high by the $\pi$-form, whereas the $\sigma$-form gives better agreement with experimental data. It seems that the hybridization type of the O -atom determines predominantly the height of rotational barrier, while the choice of the Kékulé structures are not so important.

The PCILO value of $3.0 \mathrm{kcal} / \mathrm{mol}$ with the $\sigma-(\mathrm{A}+\mathrm{B} / 2)$-form is similar to those calculated by CNDO/2 method using standard geometry (Table 2). The PCILO value of $4.94 \mathrm{kcal} / \mathrm{mol}$ with the $\pi-(\mathrm{A}+\mathrm{B} / 2)$-form is comparable with the $a b$ initio STO-3G results using standard geometry, but are still not in good agreement with the experimental values.
f) Anisole. The internal rotation of the methoxy group around the two axes $\tau_{1}(\mathrm{CCOC})$ and $\tau_{2}\left(\mathrm{COCH}^{*}\right)$ seems to be due to a balance between the steric interaction of the methyl group and the ortho-H-atoms of the phenyl system which


Kékulé A

Kékulé B


Fig. 11. Geometry of $O$-atom in anisole
favors a non-planar conformation, and the $\pi$-overlap which favors a planar heavy atom skeleton (Fig. 10). By experimental methods the conformation and rotational barrier of the methoxy group are not definitely determined. It is only indicated that the most stable conformation is planar as reviewed in [50]. The employed geometries of the fictitious atoms are given in Figure 11. In view of geometry relaxation, the bond angle $\Varangle \mathrm{CCO}$ is taken as $124.4^{\circ}$ in the planar conformation, and as $120.0^{\circ}$ in the orthogonal one. For the dihedral angle $30^{\circ}$ and $60^{\circ}$ this bond angle is linear-interpolated.

The computational results of the rotation around $\tau_{1}$, fixed with $\tau_{2}$ at $180^{\circ}$, are given in Table 7. The planar conformation is found to be the most stable one in all cases. The $\sigma$-form has a lower zeroth-order energy in the planar ( $\Delta E_{0}^{\sigma \pi}=7.57$ $\mathrm{kcal} / \mathrm{mol}$ ) and the orthogonal conformation ( $4 E_{0}^{\sigma \pi}=7.05 \mathrm{kca} /$ mol $)$. The third-order

Table 7. PCILO energy contributions of anisole with $\tau_{2}\left(\mathrm{COCH}^{*}\right)=180^{\circ}(\mathrm{kcal} / \mathrm{mol})$

| Dihedral angle $\tau_{1}$ Hybridization of oxygen atom Kékulé structure | $\tau_{1}=0^{\circ}$ |  |  |  | $t_{1}=90^{\circ}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\sigma$ |  | $\pi$ |  |  | $\pi$ |
|  | A | B | A | B | $\mathrm{A}=\mathrm{B}$ | $\mathrm{A}=\mathrm{B}$ |
| Nuclear repulsion energy | 112533.58 | 112533.58 | 112533.58 | 112533.58 | 112329.59 | 112329.59 |
| $E_{0}$ (zeroth-order energy) | -46185.865 | -46185.821 | -46178.294 | -46178.249 | -46187.712 | -46180.662 |
| Delocalization energy | -296.302 | -295.695 | -310.967 | - 309.834 | -292.747 | -301.363 |
| Intra-bond correlation energy | - 117.228 | - 117.243 | -116.549 | -116.565 | -117.235 | - 116.552 |
| 1 nter-bond correlation energy | - 124.127 | - 123.126 | - 123.621 | - 122.647 | - 123.133 | - 122.647 |
| $E_{2}$ (second-order energy correction) | -537.657 | -536.064 | -551.138 | - 549.046 | - 533.115 | - 540.562 |
| $E_{0}+E_{2}$ | -46723.522 | -46721.885 | -46729.432 | -46727.295 | -46720.826 | -46721.225 |
| Delocal. 2 bonds correl. interaction ${ }^{\text {a }}$ ) | - 10.634 | - 10.638 | - 10.558 | - 10.561 | - 10.824 | - 10.714 |
| 1 Bond correl. 2 bond correl. interaction | 48.365 | 47.891 | 48.072 | 47.617 | 47.960 | 47.681 |
| Delocal. delocal. interaction | -48.669 | -47.899 | -44.266 | -43.583 | -47.439 | -46.332 |
| 2 Bond correl. 2 bond correl. interaction | 36.103 | 35.199 | 35.199 | 35.088 | 35.521 | 35.402 |
| $E_{3}$ (third-order energy correction) | 25.162 | 24.552 | 29.213 | 28.556 | 25.217 | 26.035 |
| $E_{0}+E_{2}+E_{3}$ | -46698.360 | -46697.333 | -46700.219 | -46698.739 | -46695.609 | -46695.189 |

[^2]energy in the planar conformation is lower for the $\pi$-form ( $\Delta E_{3}^{\alpha \pi}=1.63 \mathrm{kcal} / \mathrm{mol}$ ), but in the orthogonal conformation the $\sigma$-form ( $\Delta E_{3}^{\sigma \pi}=0.42 \mathrm{kcal} / \mathrm{mol}$ ) has a lower energy. Similar to phenol this change in the third-order energy is mainly due to changes in the delocalization energy, which is $3.25 \mathrm{kcal} / \mathrm{mol}$ for the $\sigma$-form and $9.04 \mathrm{kcal} / \mathrm{mol}$ for the $\pi$-form. The energy differences between two Kékulé structures of the same hybridization type are $\Delta E_{3 \sigma}^{\mathrm{AB}}=1.0 \mathrm{kcal} / \mathrm{mol}$ and $\Delta E_{2 \pi}^{\mathrm{AB}}=1.5 \mathrm{kcal} /$ mol in the planar conformation. In the orthogonal conformation the two Kékulé structures are degenerate.

The application of the criteria for the choice of the appropriate description may be indicated as follows,

$$
\begin{aligned}
E_{0} \text {-criterion } \rightarrow & \sigma \text {-form, Kékulé structure } \mathrm{A} \\
E_{3} \text {-criterion } \rightarrow & \text { planar: } \pi \text {-form, Kékulé structure } \mathrm{A} \\
& \text { orthogonal: } \sigma \text {-form, Kékulé structure } \mathrm{A}=\mathrm{B}
\end{aligned}
$$

Compared with experimental data the choice of the $\pi$-form gives a too high barrier while by the $\sigma$-form a more reasonable value is achieved. As in phenol, the height of the rotational barrier is determined mainly by the hybridization type of the O -atom, seldom by the Kékule structures. It may be suggested that the delocalization in the $\pi$-form is overestimated in the planar conformation.

These results for the $\sigma$-form with respect to the barrier to internal rotation $\tau_{1}$ ( $2.1 \mathrm{kcal} / \mathrm{mol}$ ) agree well with those obtained from $\mathrm{CNDO} / 2$ calculations using standard geometry ( 1.0 and $1.76 \mathrm{kcal} / \mathrm{mol}$ ) and MINDO/3 ( $1.28 \mathrm{kcal} / \mathrm{mol}$ ), whereas in the ab initio STO-3G studies using standard geometry the orthogonal conformation is almost degenerate in energy with the planar conformation and INDO even fails in predicting the most stable conformation (Table 2).

The conformational energy maps of anisole with respect to the torsional angles $\tau_{1}$ and $\tau_{2}$ are shown in Figure 12. As noted from the maps the rotational barrier height is smaller in magnitude due to a rotation of the methyl group, resulting from decreased steric interaction between the methyl group and the ortho-H-atoms of the phenyl system. Furthermore it is worthwhile to note that the $\sigma$ - and $\pi$-forms do not influence the rotation about $\tau_{2}$ significantly (cf. Table 8). The rotation of the methyl group with $\tau_{1}=90^{\circ}$ is expected to be symmetrical. As shown in Figure 13, the separate Kékulé structures A and B gives an unsymmetrical potential function, whereas the superposition ( $\mathrm{A}+\mathrm{B} / 2$ ) symmetrize the results. Therefore, the ( $A+B / 2$ )-criterion seems to give the most appropriate description of the rotational behaviour of anisole.

Table 8. Rotational barriers of the methyl group in anisole in $\left.\mathrm{kcal} / \mathrm{mol}\left(\tau_{2}: 180^{\circ} \rightarrow 210^{\circ}\right)^{\text {a }}\right)$

| $\tau_{1}$ | $\stackrel{\square}{\text { a }}$ |  | $\pi$ |  |
| :---: | :---: | :---: | :---: | :---: |
|  | A | B | A | B |
| $0^{\circ}$ | 3.403 | 3.301 | 3.600 | 3.201 |
| $30^{\circ}$ | 1.309 | 1.331 | 1.622 | 1.264 |
| $60^{\circ}$ | 0.977 | 0.678 | 0.888 | 0.692 |
| $90^{\circ}$ | 0.715 | 0.715 | 0.729 | 0.730 |
| ${ }^{\text {a }}$ ) $\sigma$ and $\pi$ stand for the hybridization type and A and B for the Kékule structure. |  |  |  |  |



Fig. 12. Conformational energy maps of anisole a) $\sigma(A+B / 2)$-form, $b$ ) $\pi(A+B / 2)$-form (the isoenergy curves run from 0 to $3.5 \mathrm{kcal} / \mathrm{mol}$ in $0.5 \mathrm{kcal} / \mathrm{mol}$ steps)


Fig. 13. Rotation of the methyl group in anisole with $\tau_{I}=90^{\circ}$
3. Conclusions. - The Kékulé structures and hybridization types of O -atom suggested by the application of the $E_{0^{-}}$and $E_{3}$-criteria are collected in Table 9. In phenol and anisole the hybridization type of the O -atom determines predominantly the height of barrier to internal rotation. It seems advisable to employ an $s p^{3}$-hybridized O-atom in the PCILO computations, since in phenol and anisole this hybridization type offers results comparable with experimental values. This hybridization type is also predicted from the $E_{0}$-criterion, whereas the $E_{3}$-criterion does not give a consistent picture. The higher barrier to internal rotation with the $s p^{2}$-hybridized form of phenol and anisole may arrise from the overestimation of the delocalization energy in the planar conformation. As in the case of the rotation of the methyl group in anisole, the influence of the hybridization type of O -atom during this rotation seems to be negligible, if the rotational axis is not adjacent to the aromatic system.

As summarized in Table 9, the ( $\mathrm{A}+\mathrm{B} / 2$ )-criterion generally offers a good quantitative description of internal rotations of molecules considered in the study. The $E_{0}$-criterion does not seem to be appropriate since the barriers to internal rotation in aromatic systems predominantly arrise from delocalization effects. In computations with the PCILO method the CI-part must also be considered in aromatic systems. For molecules, whose both Kékulé structures are not comparable in their third-order energy ( $\Delta E_{3}^{\mathrm{AB}}>2 \mathrm{kcal} / \mathrm{mol}$ ), it does not seem appropriate to employ the $(A+B / 2)$-scheme for the description of the whole system. The internal rotations of benzaldehyde and aniline are better described with the Kékulé structures proposed by the $E_{3}$-criterion. It is also suggested that in the strongly localized picture of PCILO the weights of two Kékulé structures in certain molecules may be different.

The general conclusion from this study is that, for the calculation of the internal rotation of a bond adjointed to an aromatic system, a criterion of differences between the two Kékulé structures ( $\Delta E_{3}^{\mathrm{AB}} \gtrless 2 \mathrm{kcal} / \mathrm{mol}$ ) must be chosen at first, and then depending on the first criterion the application of the $E_{3^{-}}$or $(\mathrm{A}+\mathrm{B} / 2)$-criterion should be determined. These findings are summarized in Figure 14 for further applications of PCILO calculations in other aromatic systems.

For the barrier to internal rotation of the groups not directly connected to an aromatic system (e.g. the methyl group in anisole), the influence of energy differences of the two Kékulé structures is not important. But attention should be

Table 9. Predictions from variational criteria for the appropriate energy value in calculations of monosubstituted benzenes Ph-X

| X | $E_{0}$-criterion | $E_{3}$-criterion | Appropriate form <br> compared with <br> experimental data |
| :--- | :--- | :--- | :--- |
| $\mathrm{CH}_{3}$ | $\mathrm{~A}=\mathrm{B}$ | opposite | $(\mathrm{A}+\mathrm{B} / 2)$ |
| $\mathrm{NO}_{2}$ | A | A | $(\mathrm{~A}+\mathrm{B} / 2)$ |
| CHO | A | B | B |
| $\mathrm{NH}_{2}$ | $\mathrm{~A} \sim \mathrm{~B}$ | A | A |
| OH | $\sigma \mathrm{A}$ | dependent upon conformation | $\sigma \mathrm{A}$ and $\sigma(\mathrm{A}+\mathrm{B} / 2)$ |
| OCH |  | $\sigma \mathrm{A}$ | dependent upon conformation |



Fig. 14. Choice of energy criteria in treating substituted aromatic systems by the PCILO method
paid to the possibility of differences in the intramolecular interaction of atoms far from the aromatic system with those adjointed to it.

With the above mentioned remedies the PCILO method provides quite a good result in determining the barrier to internal rotations. Except for aniline, the PCILO barriers are in general smaller in magnitude compared to experimental data. Independent on the choice of Kékulé structure and hybridization type of O-atom, the most stable conformation of molecules considered in this study is predicted well by the PCILO method, whereas other semiempirical methods are shown to be not reliable in general in predicting the most stable conformation of substituted benzenes (cf. Table 2).

## REFERENCES

[1] H. Weinstein. Int. J. Quant. Chem., Quant. Biol. Symp. 2, 59 (1975).
[2] J.J. Kaufman, Int. J. Quant. Chem., Quant. Biol. Symp. 4, 375 (1977).
[3] J. B. Green, C. L. Johnson \& S. Kang, Ann. Rev. Pharmacol. 14, 319 (1974).
[4] A. Pullman, Fortschr. Chem. Forsch. 31. 45 (1972).
[5] R. E. Cristoffersen, in 'Quantum Mechanics of Molecular Conformation', B. Pullman (Ed.), J. Wiley, New York 1979.
[6] B. Pullman, Adv. Quant. Chem. 10, 251 (1977).
[7] S. Diner, J. P. Malrieu \& P. Claverie, Theoret. Chim. Acta I3, I (1969); J. P. Malrieu, P. Claverie \& S. Diner, Theoret. Chim. Acta 13, 18 (1969).
[8] J.P. Malrieu, in 'Semiempirical Methods of Electronic Structure Calculation', G.A. Segal (Ed.), Part A, Plenum Press, New York 1977, p. 69.
[9] M. Martin, R. Carbo, C. Petrongolo \& J. Tomasi, J. Am. Chem. Soc. 97, 1338 (1975).
[10] R. Cetina, M. Rubio \& O. A. Novaro, Theoret. Chim. Acta 32, 81 (1973).
[11] C. Beguin \& E. Gout-Mallaret, J. Fluor. Chem. 8, 279 (1976).
[12] Th. Weller, H.J. Hofmann \& D. Heidrich, Z. Chem. 16, 404 (1976).
[13] P. Palla, C. Petrongolo \& J. Tomasi, J. Phys. Chem. 84, 435 (1980).
[14] J. Berger \& H. Perrin, J. Mol. Struct., Theochem. 76, 375 (1981).
[15] V.A. Zubkov, T. M. Birshtein \& I. S. Milevskaya, J. Mol. Struct. 27, 139 (1975).
[16] 'Quantum Chemistry Program Exchange', Bloomington, Indiana, USA.
[17] J.A. Pople \& M. Gordon, J. Am. Chem. Soc. 89, 4253 (1967).
[18] D. G. Lister, J. K. Tyler, J. H. Hog \& N. W. Larsen, J. Mol. Struct. 23, 253 (1974).
[19] T. Pedersen, N. W. Larsen \& L. Nygaard. J. Mol. Struct. 4, 59 (1969).
[20] M. Colapietro, Acta Cryst. B34, 3277 (1978).
[21] G. E. Campagnero \& J. L. Wood, J. Mol. Struct. 6, 117 (1970).
[22] M. Colapietro, Acta Cryšt. B33, 2240 (1977).
[23] W.A. Kreiner, H.D. Rudolph \& B.T. Tan, Mol. Spectrosc. 48, 86 (1973).
[24] H.D. Rudolph, H. Preizler, A. Gaeschke \& P. Wendling, Z. Naturf. 22 A, 940 (1967).
[25] W.J. Hehre, L. Radom \& J. A. Pople, J. Am. Chem. Soc. 94, 1496 (1972).
[26] J. H. Hog, L. Nygaard \& G. O. Sorensen, J. Mol. Struct. 7, 111 (1970).
[27] T. Correll, N. W. Larsen \& T. Pedersen, J. Mol. Struct. 65, 43 (1980).
[28] L. P. Davis \& R. M. Guidsy, Aust. J. Chem. 32, 1369 (1979).
[29] J. Trotter, Tetrahedron 8, 13 (1960).
[30] W. G. Fately, R. K. Harris, F.A. Miller \& R.E. Witkowski, Spectrochim. Acta 21, 231 (1965).
[31] R. K. Kakar, E. A. Rinehardt, C. R. Quade \& T. Kojima, J. Chem. Phys. 52, 3803 (1970).
[32] O. Gropen \& H. M. Seip, Chem. Phys. Lett. 11, 446 (1971).
[33] R. Benassi, L. Schevetti \& F. Taddei, J. Chem. Soc. Perkin 2, 545 (1979).
[34] D. Perahia \& A. Pullman. Chem. Phys. Lett. 19, 73 (1973).
[35] M. Quack \& M. Stockburger, J. Mol. Spectrosc. 43, 87 (1972).
[36] J. C. Brand, D. R. Williams \& T. J. Cock, J. Mol. Spectrosc. 20, 359 (1966).
[37] R.A. Kydd \& P.J. Krueger, Chem. Phys. Lett. 49, 539 (1977).
[38] J. C. Evans, Spectrochim. Acta 16, 428 (1960).
[39] M. K. Eberhardt \& G. Chuchani, J. Org. Chem. 37, 3649 (1972).
[40] C. C. Strametz \& H. H. Schmidtke, Theoret. Chim. Acta 42, 13 (1976).
[41] P. Scharfenberg, Z. Chem. 19, 198 (1979).
[42] A. Wolf, U. Voets \& H. H. Schmidtke, Theoret, Chim. Acta 54, 229 (1980).
[43] J. D. Dill \& P. R. Schleyer, Tetrahedron Lett. 33, 2857 (1975).
[44] H. Forest \& B. P. Dailey, J. Chem. Phys. 45, 1736 (1966).
[45] W.G. Fateley, F.A. Miller \& R.E. Witkowski, Technical Report, AFML-RR-66-408, Wright Patterson Air Force Base, Ohio, cited in W.J. Orville-Thomas (Ed.), Internal Rotation in Molecules, J. Wiley, New York 1974, p. 275.
[46] H. D. Bist \& D. R. Williams, Bull. Am. Phys. Soc. Il, 826 (1966).
[47] L. Radom, W.J. Hehre, J.A. Pople, G. L. Charlson \& W. G. Fately, J. Chem. Soc. Chem. Commun. 1972, 308.
[48] S. A. Kudchadker, R. M. Hedges \& B.J. Zwalinski, J. Mol. Struct. 43, 259 (1978).
[49] M. K. Eberhardt \& G. Chuchani, J. Org. Chem. 37, 3654 (1972).
[50] G. M. Anderson, P. A. Kollman, L. N. Domelsmith \& K.N. Honk, J. Am. Chem. Soc. I01, 2344 (1979).
[51] J. Catalan \& M. Yanez, J. Am. Chem. Soc. IOI, 3490 (1979).
[52] G. Allen \& S. Fewster, in 'Internal Rotation in Molecules', W.J. Orville-Thomas (Ed.), J. Wiley, New York 1974, p. 255.
[53] D. G. Lister, J. Mol. Struct. 68, 33 (1980).
[54] E. Helgstrand, Acta Chem. Scand. 24, 3687 (1970).
[55] O. Hofer, Monatshefte f. Chem. 109, 405 (1978).
[56] T. Matsushita, Y. Osamusa \& N. Misawa, Bull. Chem. Soc. Japan 52, 2521 (1978).
[57] J. Goulon, D. Canet, M. Evans \& G.J. Davies, Mol. Phys. 30, 973 (1975).
[58] P. Diehl, H. Huber, A. C. Kunwar \& M. Reinhold, Org. Magn. Res. 9, 374 (1977).
[59] W. M. M. Bovee \& J. Smidt, Mol. Phys. 28, 1617 (1974).
[60] D. G. Lister, J.N. Mac Donald \& N. L. Owen, 'Internal Rotation and Inversion', Academic Press (1978).
[61] A. L. McClellan, 'Tables of Experimental Dipole Moment', Vol. 2, Rahara Enterprises 1974.
[62] J.P. Daudey \& J.P. Malrieu, in 'Localization and Delocalization in Quantum Chemistry', Vol. 2, O. Chalvet et al. (Ed.), D. Reidel, Dordrecht 1976, p. 175.
[63] J. Langlet, J.P. Daudey \& J. P. Malrieu, Chem. Phys. Lett. 45, 48 (1977).
[64] Th. Weller \& R. Lochmann, Acta Chim. Acad. Sci. Hung. 98, 297 (1978).


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[^1]:    ${ }^{\text {a }}$ Delocal. stands for 'delocalization' and correl. for 'correlation'.

[^2]:    ${ }^{\text {a }}$ Delocal. stands for 'delocalization' and correl. for 'correlation'.

